

Lectures on Vertex Algebras

Why vertex algebras:

- powerful computational tool in rep. theory of inf. dim Lie algebras
- Monstrous Moonshine: connection between the Monster group and the modular function j
- Zhu's Theorem on modular invariance of the space of characters of rational VOAs
- Important part of string theory (CFT): describes correlation functions

1. Basic module for affine Lie algebra \widehat{sl}_2

$$\widehat{sl}_2 = sl_2 \otimes \mathbb{C}[t, t^{-1}] \oplus \mathbb{C}K \oplus \mathbb{C}d, \quad K \text{-central}$$

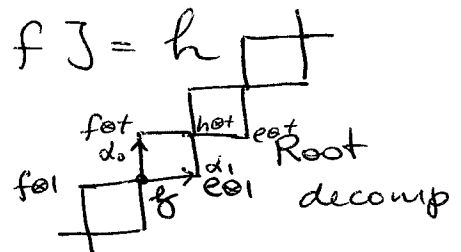
$$[xt^n, yt^s] = [x, y]t^{n+s} + n\delta_{n,-s} (x|y) K$$

$$[d, xt^n] = nxt^n$$

Basis of sl_2 : $e = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$ $h = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$, $f = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$

$$[h, e] = 2e, \quad [h, f] = -2f, \quad [e, f] = h$$

$$(e|f) = 1, \quad (h|h) = 2$$



Cartan subalgebra of \widehat{sl}_2 : $\mathfrak{h} = \langle h \otimes 1, K, d \rangle$

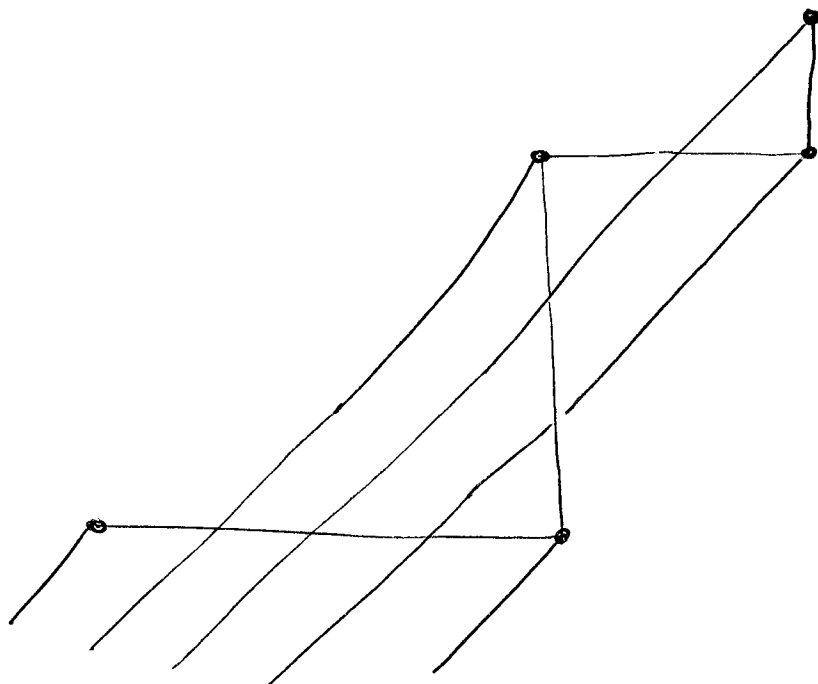
Def of a representation/module

Basic module $L(\Lambda)$ = - irreducible

highest weight module

= ... some non-trivial combinatorics ...

$$= \sum_{w \in \frac{W}{2}} e^{w(\Lambda)} \times \prod_{m \geq 1} (1 - e^{-m\alpha_0 - m\alpha_1})^{-1}$$



$\hat{\mathfrak{sl}}_2$ has an infinite-dim. Heisenberg subalgebra

$$\mathfrak{h} \otimes \mathbb{C}[t, t^{-1}] \oplus \mathbb{C}K \quad [ht^n, ht^s] = 2n\delta_{n,-s} \cdot K$$

All irreducible highest weight modules

with $\rho(K) \neq 0$ have the same structure

and their characters are

$$e^M \cdot \prod_{m \geq 1} (1 - e^{-m\delta})^{-1}, \quad \delta = \alpha_0 + \alpha_1$$

More explicitly, such a module is isomorphic

to $\mathbb{C}[x_1, x_2, x_3, \dots]$

$$\rho(ht^{-n}) = nx_n$$

$$\rho(ht^n) = 2 \frac{\partial}{\partial x_n}$$

$$\rho(k) = \text{Id}$$

$$\rho(h \otimes 1) = \alpha \cdot \text{Id}$$

Idea: coordinatize $L(\Lambda)$:

$$L(\Lambda) = \mathbb{C}[q, q^{-1}] \otimes \mathbb{C}[x_1, x_2, x_3, \dots]$$

$$\text{char}(q^n) = \Lambda - k^2 \alpha_0 - k(k-1) \alpha_1$$

$$\rho(h)q^n = 2nq^n \Rightarrow \rho(h) = 2q \frac{\partial}{\partial q}$$

We know how Heisenberg acts on $L(\Lambda)$

How to obtain the action of the rest of the $\widehat{\mathfrak{sl}}_2$?

Introduce formal power series

$$e(z) = \sum_{j \in \mathbb{Z}} et^j \cdot z^{-j-1}, \quad f(z) = \sum_{j \in \mathbb{Z}} ft^j \cdot z^{-j-1}$$

$$[ht^n, e(z)] = \sum_{j \in \mathbb{Z}} [ht^n, et^j] z^{-j-1}$$

$$= \sum_{j \in \mathbb{Z}} 2e^j t^{j+n} \cdot z^{-j-1} = 2z^n \sum_j et^{j+n} \cdot z^{-j-n-1}$$

$$[ht^n, f(z)] = -2z^n f(z) \quad \left. \begin{array}{l} = 2z^n e(z) \\ \text{achieve} \\ \text{diagonalization} \end{array} \right\}$$

Want to determine $E(z) = p(e(z))$

Think that $E(z)$ is a differential operator that is made up of operators of differentiations and multiplications by x_i .

$$\left[2 \frac{\partial}{\partial x_n}, E(z) \right] = 2 z^n E(z)$$

$\frac{\partial}{\partial x_n}$ commutes with $x_k, k \neq n$
and all $\frac{\partial}{\partial x_j}$

this shows how $E(z)$ depends on x_n

$$\left. \begin{aligned} \left[\frac{\partial}{\partial x}, f(x) \right] &= f'(x) \\ \left[\frac{\partial}{\partial x}, f(x) \right] &= \lambda f(x) \end{aligned} \right\} \Rightarrow f(x) = C \exp(\lambda x)$$

$$\Rightarrow E(z) \sim \exp(x_n z^n)$$

$$\Rightarrow E(z) \sim \exp\left(\sum_{n=1}^{\infty} x_n z^n\right)$$

Likewise $\left[n x_n, E(z) \right] = 2 z^{-n} E(z)$

shows how $E(z)$ depends on $\frac{\partial}{\partial x_n}$

$$E(z) \sim \exp\left(-\sum_{n=1}^{\infty} \frac{2z^{-n}}{n} \frac{\partial}{\partial x_n}\right)$$

$$\Rightarrow E(z) \sim q \exp\left(\sum_{n=1}^{\infty} x_n z^n\right) \exp\left(-\sum_{n=1}^{\infty} \frac{2z^{-n}}{n} \frac{\partial}{\partial x_n}\right)$$

For example,

$$E(z) \cdot 1 = q \exp\left(\sum_{n=1}^{\infty} x_n z^n\right) =$$

$$= q \cdot z^0 + q x_1 \cdot z + q\left(x_2 + \frac{x_1^2}{2}\right) z^2 + \dots$$

$$\rho(et^{-1}) \cdot 1 = q \quad \rho(et^n) \cdot 1 = 0 \text{ for } n \geq 0$$

$$\rho(et^{-2}) \cdot 1 = q x_1$$

$$\rho(et^{-3}) \cdot 1 = q\left(x_2 + \frac{x_1^2}{2}\right) \text{ etc.}$$

Does not quite work when applied to other powers of q :

should be $\rho(e \otimes 1) q^{-1} \neq 0$

~~$\rho(e \otimes 1) q^{-1} \neq 0$~~

$\rho(et) q^{-1} \neq 0$

Correct formula:

$$E(z) = q z^{2q \frac{\partial}{\partial q}} \cdot \exp\left(\sum_{n=1}^{\infty} x_n z^n\right) \exp\left(-\sum_{n=1}^{\infty} \frac{2z^{-n}}{n} \frac{\partial}{\partial x_n}\right)$$

$$F(z) = q^{-1} z^{-2q \frac{\partial}{\partial q}} \exp\left(-\sum_{n=1}^{\infty} x_n z^n\right) \exp\left(\sum_{n=1}^{\infty} \frac{2z^{-n}}{n} \frac{\partial}{\partial x_n}\right)$$

= vertex operator =

By construction, we know that ~~correct~~

we have correct Lie brackets ~~between~~

$$[\rho(ht^n), \rho(et^s)] = 2\rho(et^{n+s})$$

$$[\rho(ht^n), \rho(ft^s)] = \dots$$

We need \checkmark to check

$$[\rho(et^n), \rho(ft^s)] = \frac{1}{2} 2\rho(ht^{n+s}) + n\delta_{n,-s} \cdot \rho(K)$$

$$[\rho(et^n), \rho(et^s)] = 0$$

$$[e(z_1), f(z_2)] = \sum_{ij \in \mathbb{Z}} [et^i, ft^j] z_1^{-i-1} z_2^{-j-1}$$

$$= \sum_{ij} ht^{i+j} z_1^{-i-1} z_2^{-j-1} + \sum_i Ki z_1^{-i-1} z_2^{+i-1}$$

$$= \sum_{\substack{j \\ i+j}} ht^{i+j} z_2^{-i-j-1} \times \sum_i z_1^{-i-1} z_2^i + K \frac{\partial}{\partial z_2} \sum_i z_1^{-i-1} z_2^i$$
$$= h(z_2) \cdot z_1^{-1} \delta\left(\frac{z_2}{z_1}\right) + K \cdot z_1^{-1} \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

δ -function: $\delta(z) = \sum_{j \in \mathbb{Z}} z^j$

Properties of δ -function

$$z^n \delta(z) = \delta(z)$$

$$P(z) \delta(z) = P(1) \delta(z)$$

$$\frac{\partial}{\partial z} P'(z) \delta(z) + P(z) \delta'(z) = P(1) \delta'(z)$$

$$\Rightarrow P(z) \delta'(z) = P(1) \delta'(z) - P'(1) \delta(z)$$

$$\left(\frac{z_2}{z_1}\right) \delta\left(\frac{z_2}{z_1}\right) = \delta\left(\frac{z_2}{z_1}\right)$$

$$z_2 \delta\left(\frac{z_2}{z_1}\right) = z_1 \delta\left(\frac{z_2}{z_1}\right) \Rightarrow (z_2 - z_1) \delta\left(\frac{z_2}{z_1}\right) = 0$$

$$P(z_1, z_2) \delta\left(\frac{z_2}{z_1}\right) = P(z_2, z_2) \delta\left(\frac{z_2}{z_1}\right)$$

$$(7) = P(z_1, z_1) \delta\left(\frac{z_2}{z_1}\right)$$

$$z_2 \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right) = -z_1 \frac{\partial}{\partial z_1} \delta\left(\frac{z_2}{z_1}\right)$$

$$\frac{\partial P}{\partial z_1}(z_1, z_2) \delta\left(\frac{z_2}{z_1}\right) + P(z_1, z_2) \frac{\partial}{\partial z_1} \delta\left(\frac{z_2}{z_1}\right) = P(z_2, z_2) \frac{\partial}{\partial z_1} \delta\left(\frac{z_2}{z_1}\right)$$

$$\frac{\partial P}{\partial z_1}(z_2, z_2) \delta\left(\frac{z_2}{z_1}\right) = P(z_1, z_2) \cdot \left(\frac{z_2}{z_1}\right) \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

$$= -P(z_2, z_2) \left(\frac{z_2}{z_1}\right) \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

divide by z_2

$$\Rightarrow P(z_1, z_2) z_1^{-1} \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

$$= \frac{\partial P}{\partial z_1}(z_2, z_2) z_1^{-1} \delta\left(\frac{z_2}{z_1}\right) + P(z_2, z_2) z_1^{-1} \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

$$[E(z_1), F(z_2)] = E(z_1)F(z_2) - F(z_2)E(z_1)$$

$$= E(z_1)F(z_2) = q z_1^{2q} \frac{\partial}{\partial q} \exp\left(\sum_{n=1}^{\infty} x_n z_1^n\right) \exp\left(-\sum_{n=1}^{\infty} \frac{2z_1^n}{n} \frac{\partial}{\partial x_n}\right)$$

$$\times q^{-1} z_2^{-2q} \frac{\partial}{\partial q} \exp\left(-\sum_{n=1}^{\infty} x_n z_2^n\right) \exp\left(+\sum_{n=1}^{\infty} \frac{2z_2^{-n}}{n} \frac{\partial}{\partial x_n}\right)$$

$$= \left(\frac{z_1}{z_2}\right)^{2q} \frac{\partial}{\partial q} \cdot z_1^{-2} \exp\left(\sum_{n=1}^{\infty} x_n (z_1^n - z_2^n)\right) \exp\left(-\sum_{n=1}^{\infty} \frac{2(z_1^{-n} - z_2^{-n})}{n} \frac{\partial}{\partial x_n}\right)$$

$$\times \exp\left(+\sum_{n=1}^{\infty} \frac{2}{n} \left(\frac{z_2}{z_1}\right)^n\right)$$

$$\exp\left(+2 \sum_{n=1}^{\infty} \frac{x^n}{n}\right) = \exp(-2 \ln(1-x)) = \frac{1}{(1-x)^2}$$

$$= \frac{d}{dx} \frac{1}{1-x} = 1 + 2x + 3x^2 + \dots$$

$$z_1^{-2} \exp\left(\sum_{n=1}^{\infty} \frac{2}{n} \left(\frac{z_2}{z_1}\right)^n\right) = z_1^{-1} \left(\frac{1}{z_1} + 2 \frac{z_2}{z_1^2} + 3 \frac{z_2^2}{z_1^3} + \dots\right)$$

Combining with $F(z_2)E(z_1)$ we get

$$\left(\frac{z_1}{z_2}\right)^{2g} \frac{\partial}{\partial g} \exp\left(\sum_{n=1}^{\infty} \alpha_n (z_1^n - z_2^n)\right) \exp\left(-\sum_{n=1}^{\infty} \frac{2(z_1^{-n} - z_2^{-n})}{n} \frac{\partial}{\partial x_n}\right) \\ \times z_1^{-1} \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right) =$$

$$= \frac{\partial P}{\partial z_1}(z_2, z_2) \cdot z_1^{-1} \delta\left(\frac{z_2}{z_1}\right) + P(z_2, z_2) z_1^{-1} \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

$$P(z_2, z_2) = \text{Id}$$

$$\frac{\partial P}{\partial z_1}(z_2, z_2) = 2g \frac{\partial}{\partial g} \cdot z_2^{-1} + \sum_{n=1}^{\infty} n \alpha_n z_2^{+n-1} \\ + \sum_{n=1}^{\infty} 2 \frac{\partial}{\partial x_n} z_2^{-n-1}$$

$$[E(z_1), F(z_2)] = H(z_2) \cdot z_1^{-1} \delta\left(\frac{z_2}{z_1}\right) + \rho(K) z_1^{-1} \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

QED

Exercise $[E(z_1), E(z_2)] = 0$

Sugawara construction

Consider the Fock space $F = \mathbb{C}[x_1, x_2, x_3, \dots]$

with the action of the Heisenberg Lie algebra

$$h_n \mapsto \frac{\partial}{\partial x_n} \quad h_0 \mapsto 0$$

$$[h_n, h_s] = n\delta_{n,-s} K$$

$$h_{-n} \mapsto nx_n \quad K \mapsto \text{id}$$

Form the Heisenberg field

$$h(z) = \sum_{j \in \mathbb{Z}} h_j z^{-j-1}$$

$$\text{Exercise: } [h(z_1), h(z_2)] = K \cdot z_1^{-1} \frac{\partial}{\partial z_2} \delta\left(\frac{z_2}{z_1}\right)$$

Note: while $h(z_1)h(z_2)$ is well-defined,

$h(z) \cdot h(z)$ leads to divergence:

$$\left(\sum_{n=1}^{\infty} nx_n z^{n-1} + \sum_{n=1}^{\infty} \frac{\partial}{\partial x_n} z^{-n-1} \right) \left(\sum_{n=1}^{\infty} nx_n z^{n-1} + \sum_{n=1}^{\infty} \frac{\partial}{\partial x_n} z^{-n-1} \right)$$

apply to 1 and compute the constant term

of result:

$$\sum_{n=1}^{\infty} \frac{\partial}{\partial x_n} z^{-n-1} \cdot nx_n z^{n-1} \cdot 1 = \sum_{n=1}^{\infty} nz^{-2} = \infty$$

There is a trick to avoid divergence:

- Normally ordered product

Take a field $a(z) = \sum_{j \in \mathbb{Z}} a_j z^{-j-1}$ and split it

$$\text{into } \underbrace{a_+(z)}_{\text{non-negative pow of } z} = \sum_{j=-\infty}^{-1} a_j z^{-j-1} \quad \text{and} \quad \underbrace{a_-(z)}_{\text{neg pow of } z} = \sum_{j=0}^{\infty} a_j z^{-j-1}$$

non-negative pow of z

neg pow of z

Define

$$: a(z) b(z) : \stackrel{\text{def}}{=} a_+(z) b(z) + b(z) a_-(z)$$

Example:

$$: h(z) h(z) : = \left(\sum_{n=1}^{\infty} n \alpha_n z^{n-1} \right) h(z) + h(z) \left(\sum_{n=1}^{\infty} \frac{\partial}{\partial x_n} z^{-n-1} \right). \quad \text{- well-defined}$$

Sugawara construction:

$$\text{Expand: } \frac{1}{2} : h(z) h(z) : = \sum_{n \in \mathbb{Z}} L_n z^{-n-2}$$

Exercise Show that operators L_n on $\mathbb{C}[x_1, x_2, \dots]$ satisfy the Virasoro algebra relations:

$$[L_n, L_s] = (n-s) L_{n+s} + \delta_{n,-s} \frac{n^3-n}{12} \cdot K.$$

Fields. Locality. Operations on fields.

Let V be a vector space. Typically $V = \bigoplus_{n=0}^{\infty} V_n$

A formal power series $a(z) = \sum_{n \in \mathbb{Z}} a_n z^{-n-1}$,

$a_n \in \text{End } \bar{V}$, is called a field if

$$\forall v \in \bar{V} \quad \exists K \in \mathbb{Z} \quad \forall m \geq K \quad a_m \cdot v = 0.$$

Typically $a_m \bar{V}_n \subset \bar{V}_{n+m}$

Def. Two fields $a(z)$, $b(z)$ are called mutually local if

$$[a(z), b(w)] = \sum_{n=0}^N \frac{c_n^{[n]}(w)}{n!} z^{-1} \left(\frac{\partial}{\partial w}\right)^n \delta\left(\frac{w}{z}\right)$$

for some fields $c^{[0]}(w), \dots, c^{[N]}(w)$

Example: $E(z), F(z)$ were mutually local

$$[E(z), F(w)] = H(w) z^{-1} \frac{\partial}{\partial w} \delta\left(\frac{w}{z}\right) + K z^{-1} \delta\left(\frac{w}{z}\right)$$

Here $K = K \cdot w^0$

This is equivalent to: $(z-w)^{N+1} [a(z), b(w)] = 0$.

Note: It is non-trivial to say that

$a(z)$ is ~~not~~ local with itself.

Exercise Let $\sum_{n \in \mathbb{Z}} L_n z^{-n-2}$ be a field with

components satisfying the Virasoro relations.

Show that it is local with itself.

Operations on fields:

- $a(z) \rightarrow \frac{\partial}{\partial z} a(z)$

- Normally ordered products

For local fields $a(z), b(z)$, extract

the field $c^{[n]}(z)$ from the RHS

of the commutator formula

Borcherds: All of this can be encoded in a unifying scheme:

n -th product of fields

$$a(w)_{(n)} b(w) \stackrel{\text{def}}{=}$$

$$\text{Res}_z \left(a(z) b(w) i_w (z-w)^n - b(w) a(z) i_z (z-w)^n \right)$$

i_w - expansion in positive powers of w

i_z - " - " - " - z

$$\text{Ex } i_w (z-w)^{-1} = i_w \frac{1}{z} \cdot \frac{1}{1 - \frac{w}{z}} = z^{-1} \sum_{n=0}^{\infty} \left(\frac{w}{z} \right)^n$$

$$i_z (z-w)^{-1} = -z^{-1} \sum_{n=-\infty}^{-1} \left(\frac{w}{z} \right)^n$$

Note: for $n \geq 0$ $i_w (z-w)^n = i_z (z-w)^n = (z-w)^n$

Let us compute $a(w)_{(n)} b(w)$ for $n \geq 0$:

$$\text{Res}_z \left((z-w)^n [a(z), b(w)] \right)$$

$$= \text{Res}_z (z-w)^n \sum_{k=0}^N \frac{c^{(k)}(w)}{k!} z^{-1} \left(\frac{\partial}{\partial w} \right)^k \delta \left(\frac{w}{z} \right)$$

$$\text{For } k < n : (z-w)^n \left(\frac{\partial}{\partial w} \right)^k \delta \left(\frac{w}{z} \right) = 0$$

For $k > n$, $\left(\frac{\partial}{\partial w} \right)^k \delta \left(\frac{w}{z} \right)$ does not have terms

with $z^0, z^{-1}, \dots, z^{-k+1}$. However $(z-w)^n$ has z^0, z^1, \dots, z^n

so the residue is 0. Only the term $k=n$ contributes

and we get $a(w)_{(n)} b(w) = c^{[n]}(w)$.

Let us take $n = -1$

$$a(w)_{(-1)} b(w) = \text{Res}_z \left(a(z) i_w (z-w)^{-1} \right) b(w)$$

$$= b(w) \text{Res}_z \left(a(z) i_z (z-w)^{-1} \right)$$

$$= \text{Res}_z \left(\sum_j a_j z^{-j-1} \times z^{-1} \sum_{n=0}^{\infty} \left(\frac{w}{z} \right)^n \right) b(w)$$

$$b(w) \# \text{Res}_z \left(\sum_j a_j z^{-j-1} \times z^{-1} \sum_{n=-\infty}^{-1} \left(\frac{w}{z} \right)^n \right)$$

$$= \left(\sum_{j=-1}^{-\infty} a_j w^{-j-1} \right) b(w) + b(w) \sum_{j=0}^{\infty} a_j w^{-j-1}$$

$$= : a(w) b(w) :$$

$$\text{Res}_z \left(a(z) i_w (z-w)^{-1} \right) = a_{\#}(\#)_+$$

Differentiating k times in w

$$\text{Res} \left(a(z) i_w (z-w)^{-k-1} \right) = \frac{1}{k!} \left(\frac{\partial}{\partial w} \right)^k a(w)_+$$

Thus

$$a_{\#} a(w)_{(-k-1)} b(w) = : \left(\frac{1}{k!} \left(\frac{\partial}{\partial w} \right)^k a(w) \right) b(w) :$$

Dong's Lemma

Let the fields $a(z), b(z)$ be mutually local (including with themselves)

Then the fields $a(z), b(z), \frac{\partial}{\partial z} a(z), a(z)_{(n)} b(z)$ are mutually local.

Definition of a vertex algebra

A vertex algebra is a vector space V together with the following structure:

$\mathbb{1} \in V$ identity element = vacuum

Map $D: V \rightarrow V$ (differentiation = infinitesimal translation)

State-field correspondence map

$$Y: V \rightarrow \text{End } \bar{V} [[z, z^{-1}]]$$

For $v \in V$, write $Y(v, z) = \sum_{j \in \mathbb{Z}} v_{(j)} z^{-j-1}$

$$v_{(j)} \in \text{End } \bar{V}$$

satisfying axioms

Remark: In a usual algebra multiplication

is a map $A \rightarrow \text{End } A$

$$a \mapsto L_a, \quad L_a(b) = a \cdot b.$$

so in a vertex algebra we have inf. many products

Axioms of a vertex algebra

1. $Y(\mathbb{1}, z) = \text{Id} \cdot z^0$
2. $Y(a, z)\mathbb{1}$ contains no neg. powers of z
and $Y(a, z)\mathbb{1} \Big|_{z=0} = a$ (self-replication)
3. $Y(Da, z) = [D, Y(a, z)] = \frac{\partial}{\partial z} Y(a, z)$
4. Locality $\forall a, b \in V \exists N \in \mathbb{N}$
 $(z-w)^{N+1} [Y(a, z), Y(b, w)] = 0$

Example Let $\{a_\alpha^{[k]}(w)\}_{\alpha \in S}$ be a set of mutually local fields on a vector space U .

Add the field $\mathbb{1} = \text{Id} \cdot w^0$ and take the closure with respect to n -th products and derivatives, linear combinations. The resulting vector space V of fields is a vertex algebra.

$$D = \frac{\partial}{\partial w}$$

$$Y(a(w), z) b(w) = \sum_{n \in \mathbb{Z}} a(w)_{(n)} b(w) z^{-n-1}$$

Consequences of the axioms:

Commutator formula,

$$[Y(a, z), Y(b, w)] = \sum_{k=0}^{\infty} \frac{1}{k!} Y(a_{(k)} b, w) z^{-1} \left(\frac{\partial}{\partial w} \right)^k \delta \left(\frac{w}{z} \right)$$

in particular $a_{(k)} b = 0$ for $k \gg 0$

(so $Y(a, z)$ is a field on \bar{V})

More generally,

$$Y(a_{(n)} b, z) = Y(a, z)_{(n)} Y(b, z)$$

e.g. $Y(a_{(-1)} b, z) = : Y(a, z) Y(b, z) :$

Let us revisit the basic module for \widehat{sl}_2

$$\mathbb{C}[q, q^{-1}] \otimes \mathbb{C}[x_1, x_2, \dots]$$

Generalizing operators $E(z), F(z)$, we construct vertex operators:

$$Y(q^n, z) = q^n z^{2nq \frac{\partial}{\partial q}} \exp\left(n \sum_{j=1}^{\infty} x_j z^j\right) \exp\left(-n \sum_{j=1}^{\infty} \frac{2z^j}{j} \frac{\partial}{\partial x_j}\right)$$

We also have the field

$$H(z) = \sum_{j=1}^{\infty} j x_j z^{j-1} + 2q \frac{\partial}{\partial q} z^{-1} + \sum_{j=1}^{\infty} 2 \frac{\partial}{\partial x_j} z^{-j-1}$$

$$H(z) = Y(\ ?, z)$$

Use the self-replication axiom

$$a = Y(a, z) \mathbb{1} \Big|_{z=0} = H(z) \mathbb{1} \Big|_{z=0} = x_1$$

Then we get

$$\text{Note } (x_1)_{(-n)} = n x_n$$

$$Y(x_1 q^k, z) = :Y(x_1, z) Y(q^k, z):$$

$$=: H(z) Y(q^k, z):$$

$$Y(x_{j_1} x_{j_2} \dots x_{j_s} q^k, z) =$$

$$= Y\left(\frac{1}{j_1} (x_1)_{(-j_1)} \dots \frac{1}{j_s} (x_1)_{(-j_s)} q^k, z\right) =$$

$$= \frac{1}{j_1! \dots j_s!} : \left(\frac{\partial}{\partial z}\right)^{j_1-1} H(z) \dots : \left(\frac{\partial}{\partial z}\right)^{j_s-1} H(z) Y(q^k, z) : \dots :$$

Existence Theorem (E. Frenkel-Kac-Radul-Wang)

Suppose we have $(V, \mathbb{1}, D)$

and a collection of fields $\{a^\alpha(z)\}_{\alpha \in S}$ on V

satisfying

$$(1) \quad [D, a^\alpha(z)] = \frac{\partial}{\partial z} a^\alpha(z)$$

$$(2) \quad D\mathbb{1} = 0, \quad a^\alpha(z)\mathbb{1} \in V[[z]]$$

with $\{a^\alpha(z)\mathbb{1}\}_{\alpha \in S}$ - linearly independent

(3) $\{a^\alpha(z)\}_{\alpha \in S}$ is mutually local

$$(4) \quad V = \text{Span} \left\{ a_{(-j_1-1)}^{\alpha_1} \cdots a_{(-j_k-1)}^{\alpha_k} \mathbb{1} \right\}$$

Then \bar{V} is a vertex algebra with Υ map given by

$$\Upsilon(a_{(-j_1-1)}^{\alpha_1} \cdots a_{(-j_k-1)}^{\alpha_k} \mathbb{1}, z) =$$

$$= \frac{1}{j_1! \cdots j_k!} \cdot \left(\frac{\partial}{\partial z} \right)^{j_1} a^{\alpha_1}(z) \cdots \cdot \left(\frac{\partial}{\partial z} \right)^{j_{k-1}} a^{\alpha_{k-1}}(z) \left(\frac{\partial}{\partial z} \right)^{j_k} a^{\alpha_k}(z)$$

By existence theorem, the basic module is a vertex algebra.

This can be generalized to an arbitrary integral lattice (19)

- Two main sources of vertex algebras
- integral lattices
 - infinite-dimensional Lie algebras

Vertex Lie algebras

Let \mathcal{L} be a Lie algebra with a basis

$$\{ u(n), c(-1) \mid u \in \mathcal{U}, c \in \mathcal{C}, n \in \mathbb{Z} \}$$

where $c(-1) \in \mathbb{Z}(\mathcal{L})$.

Form the generating series

$$u(z) = \sum_{n \in \mathbb{Z}} u(n) z^{-n-1}$$

$$c(z) = c(-1) \cdot z^0$$

Let \mathcal{F} be a subspace in $\mathcal{L}[[z, z^{-1}]]$

spanned by $\{ u(z), c(z) \}$ and their derivatives of all orders.

Def. \mathcal{L} is a vertex Lie algebra if

$$\forall a, b \in \mathcal{U} \quad \exists f^{[0]}(z), \dots, f^{[N]}(z) \in \mathcal{F}$$

so that

$$[a(z), b(w)] = \sum_{k=0}^N \frac{f^{[k]}(w)}{k!} z^{-1} \left(\frac{\partial}{\partial w} \right)^k \delta\left(\frac{w}{z}\right).$$

Example: Affine Lie algebras

$$\mathfrak{g} \otimes \mathbb{C}[t, t^{-1}] \oplus \mathbb{C}K$$

$$\mathcal{U} = \text{basis of } \mathfrak{g}, \quad \mathcal{C} = \{K\}$$

$$x(n) = x \otimes t^n, \quad K = K(-1)$$

Verify that it is a vertex Lie algebra:

$$[x(z), y(w)] = \sum_{i, j \in \mathbb{Z}} [x t^i, y t^j] z^{-i-1} w^{-j-1}$$

$$= \sum_{i, j} [x, y] t^{i+j} w^{-i-j-1} \cdot z^{-i-1} w^i$$

$$+ \sum_i i (x|y) \cdot K z^{-i-1} w^{i-1}$$

$$= [x, y](z) \cdot z^{-1} \delta\left(\frac{w}{z}\right) + (x|y) K(w) z^{-1} \frac{\partial}{\partial w} \delta\left(\frac{w}{z}\right).$$

Theorem (Dong-Li-Mason)

Let \mathcal{L} be a vertex Lie algebra.

Construct the subspaces:

$$\mathcal{L}_+ = \text{Span} \{ u(n) \mid n \geq 0, u \in \mathcal{U} \} \quad \begin{array}{l} \text{- coeff at} \\ \text{neg. pow of } z \end{array}$$

$$\mathcal{L}_- = \text{Span} \{ u(n), c(-1) \mid n < 0, u \in \mathcal{U}, c \in \mathcal{C} \}$$

Then

(1) Both \mathcal{L}_+ and \mathcal{L}_- are subalgebras in \mathcal{L} .

(2) Let $\mathbb{C}\mathbb{1}$ be a trivial 1-dim. repres. for \mathcal{L}_+

Construct the induced module

$$\bar{V}_{\mathcal{L}} = \text{Ind}_{\mathcal{L}_+}^{\mathcal{L}} \mathbb{C}\mathbb{1} \cong U(\mathcal{L}_-) \otimes_{\mathbb{C}} \mathbb{1}.$$

$\bar{V}_{\mathcal{L}}$ has a structure of a vertex algebra

with

$$\begin{aligned} Y(u_1(-j_1-1) \dots u_k(-j_k-1)\mathbb{1}, z) &= \\ &= \frac{1}{j_1! \dots j_k!} : \left(\frac{\partial}{\partial z} \right)^{j_1} u_1(z) \dots : \left(\left(\frac{\partial}{\partial z} \right)^{j_{k-1}} u_{k-1}(z) \right) \left(\frac{\partial}{\partial z} \right)^{j_k} u_k(z) : \end{aligned}$$

(3) Let χ be a central character $\chi: \mathbb{C} \rightarrow \mathbb{C}$

Then $\bar{V}_{\mathcal{L}}(\chi) = \bar{V}_{\mathcal{L}} / \langle \mathbb{C}(-1)\mathbb{1} - \chi(\mathbb{C})\mathbb{1} \rangle$ is a vertex algebra

(4) Let P be a maximal \mathcal{L} -submodule in $\bar{V}_{\mathcal{L}}(\chi)$.

Then $L(\chi) = \bar{V}_{\mathcal{L}}(\chi) / P$ is a simple vertex algebra

(5) Any bounded module for \mathcal{L} is a module for the vertex algebra $\bar{V}_{\mathcal{L}}$.

Example Affine Lie algebras

$$\mathcal{L}_+ = \mathfrak{g} \otimes \mathbb{C}[t]$$

$$\mathcal{L}_- = \mathfrak{g} \otimes t^{-1}\mathbb{C}[t^{-1}] \oplus \mathbb{C}K$$

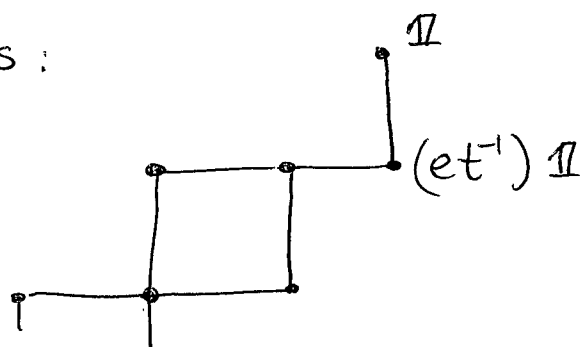
$$V_{\mathcal{L}} = U(\mathfrak{g} \otimes t^{-1}\mathbb{C}[t^{-1}] \oplus \mathbb{C}K) \otimes \mathbb{1}$$

$$V_{\mathcal{L}}(c) \cong U(\mathfrak{g} \otimes t^{-1}\mathbb{C}[t^{-1}]) \otimes \mathbb{1}, \quad K \mapsto c \cdot \text{id}$$

Let $c = 1$, $\mathfrak{g} = \mathfrak{sl}_2$ and consider the simple quotient

$L(c)$. Then $L(c)$ is the basic module.

Applications:



$$\text{In } L(1) \quad (et^{-1})^2 \mathbb{1} = 0$$

$$Y((et^{-1})^2 \mathbb{1}, z) = Y(e(-1)e(-1)\mathbb{1}, z)$$

$$= : e(z) e(z) :$$

Then $: e(z) e(z) ; = 0$ on $L(1)$.

Virasoro vertex algebra

$$[L(n), L(s)] = (s-n) L(n+s) + \delta_{n,-s} \frac{n^3-n}{12} K$$

$$\mathcal{U} = \{\omega\}, \quad \mathcal{C} = \{K\}$$

$$\omega(n) = L(n-1)$$

$$\omega(z) = \sum_{n \in \mathbb{Z}} \omega(n) z^{-n-1} = \sum_{n \in \mathbb{Z}} L(n) z^{-n-2}$$

Exercise

$$[\omega(z), \omega(w)] = \frac{\partial}{\partial w} \omega(w) z^{-1} \delta\left(\frac{w}{z}\right)$$

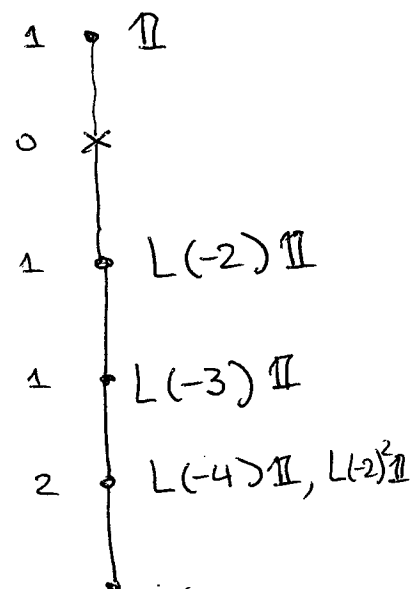
$$+ 2\omega(w) z^{-1} \frac{\partial}{\partial w} \delta\left(\frac{w}{z}\right)$$

$$+ \frac{1}{12} K(w) z^{-1} \left(\frac{\partial}{\partial w}\right)^3 \delta\left(\frac{w}{z}\right)$$

$$\mathcal{L}_- = \text{Span} \{L(n) \mid n \leq -2, K\}$$

$$\mathcal{L}_+ = \text{Span} \{L(n) \mid n \geq -1\}$$

$$V_{\mathbb{Z} \text{ vir}}(c)_{\mathbb{C}} = U(\mathcal{L}_{\leq -2}) \otimes \mathbb{1}$$



Sugawara construction revisited

$$\mathfrak{h} = \langle h(n), K \rangle$$

$$[h(n), h(s)] = n \delta_{n,-s} K$$

$$\mathcal{U} = \{h\}, \mathcal{C} = \{K\}$$

$$[h(z), h(w)] = K(w) z^{-1} \frac{\partial}{\partial w} \delta\left(\frac{w}{z}\right).$$

$$\bar{V}_{\mathfrak{h}, \mathcal{C}}(1) = \bar{U}(\mathfrak{h}_{\leq -1}) \otimes \mathbb{1}$$

$$= \mathbb{C}[x_1, x_2, x_3, \dots]$$

$$h(n) = \frac{\partial}{\partial x_n} \quad \text{H(z)} = Y(x_1, z)$$

$$h(-n) = n x_n$$

Sugawara: $\omega(z) = \frac{1}{2} : H(z) H(z) :$

is a Virasoro field.

Note: $\omega = \frac{1}{2} x_1^2$

We need to prove

$$[\omega(z), \omega(w)] = \dots$$

Use the commutator formula

$$[Y(\omega, z), Y(\omega, w)] = \sum_{k=0}^N \frac{1}{k!} Y(\omega_{(k)} \omega, w) z^{-1} \frac{\partial^k}{\partial w^k} \delta$$

We need to compute $\omega_{(n)}\omega$ for $n \geq 0$

$n=0$: $\omega_{(0)}$ is z^{-1} -term in $\frac{1}{2} : H(z)H(z)$:

apply to $\omega = \frac{1}{2} x_1^2$

$$\frac{1}{4} \left(\sum_{n=1}^{\infty} n x_n z^{n-1} \right) \left(\sum_{n=1}^{\infty} n x_n z^{n-1} + \sum_{n=1}^{\infty} \frac{\partial}{\partial x_n} z^{-n-1} \right) x_1^2$$

$$+ \frac{1}{4} \left(\sum_{n=1}^{\infty} n x_n z^{n-1} + \sum_{n=1}^{\infty} \frac{\partial}{\partial x_n} z^{-n-1} \right) \left(\sum_{n=1}^{\infty} \frac{\partial}{\partial x_n} z^{-n-1} \right) x_1^2$$

$$n=0 : z^{-1} : \omega_{(0)}\omega = \frac{1}{4} 2x_2 \frac{\partial}{\partial x_1} \cdot x_1^2 + \frac{1}{4} 2x_2 \frac{\partial}{\partial x_1} x_1^2$$

$$= 2x_1 x_2 = h(-2)h(-1)\mathbb{1}.$$

$$n=1, z^{-2} : \omega_{(1)}\omega = \frac{1}{4} x_1 \frac{\partial}{\partial x_1} \cdot x_1^2 + \frac{1}{4} x_1 \frac{\partial}{\partial x_1} \cdot x_1^2$$

$$= x_1^2 = 2\omega$$

$$n=2 \quad \omega_{(2)}\omega = 0$$

$$n=3 \quad \omega_{(3)}\omega = \frac{1}{4} \left(\frac{\partial}{\partial x_1} \right)^2 x_1^2 = \frac{1}{2} \mathbb{1}$$

$$\omega_{(n)}\omega = 0 \quad \text{for } n > 3.$$

$$[\omega(z), \omega(w)] = Y(\omega_{(0)}\omega, w) z^{-1} \delta\left(\frac{w}{z}\right)$$

$$+ Y(\omega_{(1)}\omega, w) z^{-1} \frac{\partial}{\partial w} \delta\left(\frac{w}{z}\right)$$

$$+ \frac{1}{3!} Y(\omega_{(3)}\omega, w) z^{-1} \left(\frac{\partial}{\partial w} \right)^3 \delta\left(\frac{w}{z}\right)$$

$$Y(\omega(\omega), \omega, z) = Y(h(-2)h(-1)\mathbb{I}z)$$

$$= \vdots \left(\frac{\partial}{\partial z} H(z) \right) H(z) \vdots = \frac{\partial}{\partial z} \omega(z)$$

$$[\omega(z), \omega(w)] = \frac{\partial}{\partial w} \omega(w) z^{-1} \delta\left(\frac{w}{z}\right)$$

$$+ 2 \omega(w) z^{-1} \frac{\partial}{\partial w} \delta\left(\frac{w}{z}\right)$$

$$+ \frac{1}{12} K(w) z^{-1} \left(\frac{\partial}{\partial w}\right)^3 \delta\left(\frac{w}{z}\right).$$